#### **Scaling Limits**

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- We will look at some examples.

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#### Variational form

### Variational formMinimize

 $\sum [u(x) - u(y)]^2$  $x,y:x\simeq y$ 

over  $u: u = g \in G^c$ 

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$$\int_{G} |\nabla u|^2 dx$$

over u : u = g on  $\partial G$ .

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E[g(x<sub>τ</sub>)]

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 $|L_h \to L = \nabla \cdot \bar{a} \nabla|$ 

$$u_t = L_h u; \quad u(0, x) = f(x)$$

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#### $L_h u = f \text{ for } x \in G; \ u(y) = g(y) \text{ for } y \in \partial G$

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$$\langle \xi, \bar{a}\xi \rangle = \inf_{w} \int_{T^d} \langle (\xi - \nabla w), a(x)(\xi - \nabla w) \rangle dx$$

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In 
$$d = 1$$
,  $\bar{a} = [\int_0^1 \frac{1}{a(x)} dx]^{-1}$ 

Scaling Limits -p.7/29
#### **3.** Random Medium. (Stationary and ergodic)

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$$L_{h,\omega} = \nabla \cdot a(T_{\frac{x}{h}}\omega)\nabla$$

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$$\int w^* dP = 0, D_i w_j^* = D_j w_i^*$$

$$L_h = \sum_{i,j} a_{i,j} \left(\frac{x}{h}\right) D_i D_j$$

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•  $\int \phi dx = 1, \phi > 0.$  Periodic

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In d = 1 the same answer  $\int [\frac{1}{a(x)} dx]^{-1}$ 

Scaling Limits – p. 9/29

$$L_{h,\omega} = \sum_{i,j} a_{i,j} (T_{\frac{x}{h}}\omega) D_i D_j$$

#### and

 $\bar{a} = E[a(\omega)\phi(\omega)]$ 

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It is the unique weak sense solution of

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 $\square D_i$  is well defined on  $L_2(P)$ 

Scaling Limits – p. 10/29

#### **6**. Nonlinear versions.

$$u_t + \frac{1}{2}\Delta u + H(x, \nabla u) = 0; \quad u(T, x) = f(\frac{x}{T})$$

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$$u_t + \bar{H}(\nabla u) = 0; \quad u(1, x) = f(x)$$

Scaling Limits – p. 11/29

#### • How is $\overline{H}(p)$ related to H(x, p)?

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A<sup>\*</sup><sub>b</sub>φ<sub>b</sub> = 0

How is H(p) related to H(x, p)?  $\square L(x,q)$  is the Legendre transform. • Consider  $\mathcal{A}_b = \frac{1}{2}\Delta + \langle b(x), \nabla \rangle$  on the torus.  $\blacksquare \mathcal{A}_b^* \phi_b = 0$  $\bar{H}(p) =$  $\sup_{b(\cdot)} [< p, \int b(x)\phi_b(x) > - \int L(x, b(x))\phi_b(x)]$ 

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- J.L.Lions, P.L.Lions, A.Bensoussan,
   G.C.Papanicolaou, F.Rezakhanlou, E.Kosygina, P.E.
   Suganidis & V.

#### **7**. Interacting particle systems.

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- How does it evolve?

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Start far away from equilibrium.

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**St**art far away from equilibrium.

**How does the density evolve to equilibrium?** 

 $F_J(\eta) = \frac{1}{N^d} \sum J(\frac{x}{N})\eta(x)$ 

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### **Compute** $N^k \mathcal{L} F_J$

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## Compute $N^k \mathcal{L} F_J$ k = 1 if $\sum_z zp(z) = m \neq 0$ k = 2 if $\sum_z zp(z) = 0$

$$F_J(\eta) = \frac{1}{N^d} \sum J(\frac{x}{N})\eta(x)$$

Compute N<sup>k</sup> LF<sub>J</sub>
k = 1 if ∑<sub>z</sub> zp(z) = m ≠ 0
k = 2 if ∑<sub>z</sub> zp(z) = 0
If p(z) = p(-z) it is a lot easier.

$$N^{2-d} \sum_{x,y} \eta(x)(1-\eta(y))p(y-x)[f(\eta^{x,y}) - f(\eta)]$$
  
=  $\frac{N^{2-d}}{2} \sum_{x,y} [\eta(x) - \eta(y)]p(y-x)J[(\frac{y}{N}) - J(\frac{x}{N})]$   
 $\simeq \frac{1}{2N^d} \sum_x C_{i,j}(\partial_i \partial_j f)(\frac{x}{N})$ 

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• where  $C_{i,j} = \sum_{z} z_i z_j p(z)$ 

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$$\rho_t = \frac{1}{2} \nabla \cdot C \nabla \rho - \nabla \cdot m \rho (1 - \rho)$$

The average of  $\eta(x)(1 - \eta(y))$  is replaced by its local expectation  $\rho(1 - \rho)$ .

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 $\Box C$  is not easily computable.

#### Generator

 $\mathcal{L}_{N} = N^{2} \sum p(y - x)\eta(x)(1 - \eta(y))[F(\eta^{x,y}) - F(\eta)]$ x,y

#### **Generator**

$$\mathcal{L}_{N} = N^{2} \sum_{x,y} p(y-x)\eta(x)(1-\eta(y))[F(\eta^{x,y})-F(\eta)]$$

$$F_J(\eta) = \langle J, \rho \rangle = \frac{1}{N^d} \sum_x J(\frac{x}{N}) \eta(x)$$

$$\mathcal{L}_{\mathcal{N}}\mathcal{F}_{\mathcal{J}} = N^{2-d} \sum_{x,y} p(y-x)\eta(x)(1-\eta(y))$$
$$\times \left[J(\frac{y}{N}) - J(\frac{x}{N})\right]$$

$$\mathcal{L}_{\mathcal{N}}\mathcal{F}_{\mathcal{J}} = N^{2-d} \sum_{x,y} p(y-x)\eta(x)(1-\eta(y))$$
$$\times \left[J(\frac{y}{N}) - J(\frac{x}{N})\right]$$
$$\simeq N^{1-d} \sum_{x} \nabla J(\frac{x}{N}) \cdot W(x,\eta)$$

#### • if $p(\cdot)$ is symmetric

$$W_{i}(x,\eta) = \frac{1}{2} \sum_{j} C_{i,j}[\eta(x+e_{j}) - \eta(x)]$$

Can do summation by parts.

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The ?? can be ignored

#### • $E[W_i] = 0$ in all equilibria form a Hilbert space.

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Density gradients η(x + e<sub>j</sub>) - η(x) are complementary.

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- **There is subspace that are negligible.**.
- **Codimension** *d* in the Hilbert space.
- Density gradients  $\eta(x + e_j) \eta(x)$  are complementary.
- **Done in each equilibrium**  $P_{\rho}$  with  $C_{i,j}(\rho)$

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- **There is subspace that are negligible.**.
- Codimension *d* in the Hilbert space.
- Density gradients  $\eta(x + e_j) \eta(x)$  are complementary.
- Done in each equilibrium P<sub>ρ</sub> with C<sub>i,j</sub>(ρ)
   Large Deviation theory.

#### **8**. Back to the symmetric case. Tagged Particles.

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Diffuses.  $\lambda^{-1}x(\lambda^2 t) \to B(t)$ .

# 8. Back to the symmetric case. Tagged Particles. What about the motion of a tagged particle in equilibrium at density ρ? Diffuses. λ<sup>-1</sup>x(λ<sup>2</sup>t) → B(t).

**Covariance** is  $S(\rho)$ 

$$L_t = \frac{1}{2}\nabla \cdot S(\rho(t,x))\nabla + \frac{(S(\rho(t,x)) - C)\nabla\rho}{2\rho}\nabla$$

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• Markov with generator  $L_t$ .

$$\dot{q}_i = p_i, \quad \dot{p}_i = -\sum_j (\nabla V)(x_i - x_j)$$

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#### **Conserved quantities.** $\rho$ , u, e.

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First order hyperbolic PDE for them.
Connect the ODE's to PDE

#### **Gibbs States.** Constant $\rho$ , u, e.

### Gibbs States. Constant ρ, u, e. Local Gibbs state, Slowly varying ρ, u, e

$$\begin{bmatrix} \rho_0, u_0, T_0 \end{bmatrix} \longrightarrow Local Gibbs \downarrow Euler \qquad \qquad \downarrow Liouville \begin{bmatrix} \rho_t, u_t, T_t \end{bmatrix} \longrightarrow do not match$$

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Diagram does not commute!

$$[\rho_0, u_0, T_0] \longrightarrow Local Gibbs \downarrow Euler \qquad \qquad \downarrow Liouville [\rho_t, u_t, T_t] \longrightarrow do not match$$

Diagram does not commute!

It almost does after some noisy modification.

### Work done over 25 years. Presutti, De Masi, H.T.Yau, Olla, Rezakhanlou, Quastel, Kosygina, V

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#### **Thank You**